Compensation of Fiber Dispersion with an Ideal Dispersion Component

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B.Sc (Honours) in Engineering Technology, University of Gezira, (2001)

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Computer Engineering Department

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Date: January /2014
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Compensation of Fiber Dispersion with an Ideal Dispersion Component

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Abstract

The main objective of this project is to study the effect of the fiber dispersion in single mode fiber and show how it can compensate. In order to achieve the objectives of the project, we studied and stimulated the phenomenon referred to as group-velocity dispersion (GVD) in Single-Mode Fibers and shown how dispersion of the pulse limits the distance because the receiver must recover an unstable signal. The ability of the receiver to do this varies quite a lot with the design of the receiver itself and with the communication protocol. This casing a signal degradation and make it spreads and broaden along the distant and that lead to reduced the system performance; then shown how we can control and manage this effecting. the OptiSystem Software was used to demonstrate the influence of the (GVD) on Gaussian pulse propagation, and shown how the single was broadened; then compensate of this dispersion and single broaden with Ideal dispersion components. After the most major important factor was set like bit rate equal to 40 GB/S, which correspond to bit duration of 25 ps; and using the default value of 0.5 for “width” of the Optical Gaussian Pulse Generator and the resulting FWHM of the pulse is 12.5 ps. Then calculated the dispersion length which was 2.812km and set it at the fiber properties and disable all the effects except GVD and the project calculated under these values, the obtained result shown that the pulse was broaden (the peak power decreased). Then changed the fiber length to (10 km) and calculated the project the result shown that how the signal broaden and the width of the pulse increased approximately four time; Then added the Ideal dispersion compensation component to the design and set the dispersion parameter to (-160 ps/nm) and calculated the project; the result shown that how compensation of accumulated dispersion was achieved. In this project, the properties of optical fibers are discussed. When the GVD is considered, and shown how it affected the signal and reduced the system performance and limiting the transmission distant. Compensation necessary. However, dispersion compensation introduces additional loss and complexity to systems. This leads to various system configurations in terms of transmission fibers deployed.
تعرض الفقد الذي يحدث في الإشارة نتيجة التشتت في كواكب الألياف الضوئية

هناك عبد الرحمن إبراهيم مصطفى

من خلال الدراسة

ظاهرة (مجموع سرعة التشتت – GVD) التي تؤدي إلى انتشار وتشتت الإشارة في الألياف الضوئية

أحادية النموذج، مما يؤدي إلى ضعف عملية الاتصال وذلك لأن الطرف المستقبل للنضبة يعمل على
استهداد إشارات غير مستقرة. وتختلف اجهزة الاستقبال من حيث الأمكانية للقيام بذلك حسب التصميم
وبيوترنوكولات الاتصال المستخدم في عملية الرسالة. وهذا في مجمله يؤدي إلى اضمحلال الإشارة
وتشتت وتسع عرض نطاقها أثناء عملية الرسالة مما يؤدي بدوره إلى كفاءة وفعالية
نظام الاالصال المستخدم بصورة عامة. الهدف الرئيسي هو دراسة تأثير تشتت
وانتشار الإشارة الضوئية في الألياف الضوئية أحادية المود. وتقديم طريقة يتم بها تعيث عرض الفقد
والذي يحدث في الإشارة الضوئية. من أجل تحقيق أهداف هذا المشروع تم استخدام نظام النم زجه
والمحاكاة (أوتي سيسيم) للنزعة ومقاومة النضبة وتوضيح كيف يتم توسيع نطاق الإشارة نتيجة
الانتشار. تم دراسة ونلمة كيفية التغيير عن هذا الانتشار والتوسع الذي يتم في عرض نطق
الأشارة باستخدام (مكونات التشتت المثالية) على نظام النم زجه والمحاكاة.

وتم ضبط بعض البارامترات والموهارات لعملية النم زجه والمحاكاة ومنها معدل أرسل البيانات والذي يكافي
40 قيبايت/ثانية والتي تتوافق مع فترة ارسال تعادل 25 بيكو/ثانية، مع استخدام القيمة الإفتراضية
لعرض نطاق مدند النبضات والتي تعادل (0.5)، الذي ينتج عنه نبضة بفتره ارسل تعادل 12.5
بيكو/ثانية. ومن ثم حساب طول الانتشار في ظل هذه البارامترات وقيمته 2.812 كيلومتر، ومن ثم
ضبط خصائص كليب الألياف الضوئية المستخدم في عملية النم زجه والمحاكاة على هذه القيمة وكذلك
تطلب كل العوامل المؤثر الأخرى بانتظار تأثير ال
- GVD). والنتائج المستخلصة من هذه العملية
نجد أن نطاق عرض النبضة قد تسعك مما قد إلى انخفاض قيمة قوة الزرو للفضية ، ولأنه ان
معال الانتشار والشتت يزيد مع زيادة مسافة الاتصال أي زيادة طول كليب الألياف الضوئية
المستخدم ثم ضبط خصائص كليب الألياف الضوئية على القيمة 10 كيلومتر. ووجد أن الإشارة قد
تتشتت وتستحع وان عرض النبضة قد تقاس عقبا بدقة ارجاع عرض النبضة الأصلية . واخيرا
تمت اضافة العصر الذي استخدم في معالجة الاشارة (مكونات التشتت المثالية) للتعامل مع
التشويه الذي تحتاج عملة التشتت والزائدة في عرض النبضة الذي حدث لإعلان التشتت
معال الانتشار وهو عبارة عن القيمة - 160 بيكوثنان/نانومتر ومن ثم تشغيل البرنامج ، ومن
النتائج المستخلصة لوحظ أن قد تم تسجيل الإشارة والتحكم في عمليه التشتت والانتشار ومن ثم قد
تم ارجاع عرض النبضة لقيمة النهاية . مما يؤدي لسهولة التعامل معها لكافة الأساس وما
يحتاجها فعالية وكفاءة النظام. من اهم البارامترات التي يجب مراجعتها عند تصميم أنظمة الاتصالات
باستخدام الألياف الضوئية ذات السرعات العالية والتي تتطلب معدل ارسال بينات على مسافات طويلة
استخدام احدى تهتز على معالجة الإشارة للحد من الفقد والعشوائي الذي يحدث للاشارة نتيجة ظاهرة
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Chapter One
Introduction
1.1 Introduction

- Dispersion in Optical Fibers

Dispersion as general is a phenomenon in which the velocity of propagation of any electromagnetic wave is wavelength dependent; for examples [1]:

i. In optics, dispersion is the phenomenon in which the phase velocity of a wave depends on its frequency, or alternatively when the group velocity depends on the frequency.

ii. In communication, dispersion is used to describe any process by which any electromagnetic signal propagating in a physical medium is degraded because the various wave characteristics (i.e., frequencies) of the signal have different propagation velocities within the physical medium.

- Sources or Type of Dispersion

In general there are two sources or type of dispersion in single mode fiber [1]:

i. Material dispersion which comes from a frequency-dependent response of a material to waves.

ii. Waveguide dispersion occurs when the speed of a wave in a waveguide (such as an optical fiber) depends on its frequency for geometric reasons, independent of any frequency dependence of the materials from which it is constructed. More generally, "waveguide" dispersion can occur for waves propagating through any inhomogeneous structure. In general, both types of dispersion may be present, although they are not strictly additive. Their combination leads to signal degradation in optical fiber for telecommunication, because the varying delay in arrival time between different components of a signal "spread out" the signal in time. Dispersion in fibers is one of the limiting factors that determine how much data can be transported on a single fiber.
1.2 Problem Statement

When the dispersion occur the pulse of light is spread out during transmission on the fiber. A short pulse becomes longer and ultimately joins with the pulse behind, making recovery of a reliable bit stream impossible. Each pulse broadens and overlaps with its neighbors, eventually becoming indistinguishable at the receiver input.

1.3 Objectives

The primarily objectives of this research are to:

- Study the optical transmission system concepts and how its effect the communication world.
- Study the fiber optical as it’s the essential things when talking about the optical transmission and this thesis focusing on the single mode fiber.
- Study how light confined and transmitting on a fiber and what kind of problems that could faced.
- Study the effect of the Group Velocity Dispersion (GVD) on Gaussian Pulse Propagation; how it affects the single during the transmission cycle and how can minimize this.
- Study the dispersion phenomenon and how its reduce the communications systems performance.
- Study how the dispersion cause the pulse broaden and dispersive and how it can be recover and compensate by using Ideal Dispersion Component.

1.4 Methodology

OptiSystem Software is used to demonstrate the influence of the GVD on Gaussian pulse propagation. The Bill of material used to demonstrate the influence of the group (GVD) velocity dispersion on pulse propagation in optical fibers in "linear" regime as following [2]:

i. User Defined Bit Sequence Generator.
ii. Optical Gaussian pulse Generator.
iii. Optical Time Domain Visualizer.
iv. Optical Spectrum Analyzer.
v. Optical Fiber.

The most major important factors to set are [2]:

i. Bit rate.
ii. Bit duration.
iii. Value for width.
iv. Fiber length (Dispersion length).

Then calculate the project under these values, the obtain result should show that the pulse is broadened (the peak power decreases) [2].

To demonstrates the possibility of dispersion compensation with the help of Ideal Dispersion Component in OptiSystem. We use the same last circuit and components with the same values, in addition to Ideal Dispersion Component; the major parameter to demonstrate this is frequency; bandwidth and dispersion parameter.

1.5 Thesis Outline

This Dissertation contains five chapters; chapter one is an introduction about the thesis outline; when chapter two included literature review and full details about the major points related to the subject of the thesis; on other hand chapter three contain full description about the circuits design and implementation using the optiystem program; then the whole results that can obtained by running the circuits which mention at chapter three presented at chapter four results and discussion and at the last chapter five included the conclusion and recommendation.
Chapter Two
Literature Review
2.1. **Basic Concepts of Fiber Optics**

The major Goal is to communicate information and it can be summarized as flow [3]:

i. Transmitter sends signal; each signal has a detected data rate and format and signal modulates a carrier.

ii. Signal goes through a length of fiber; during this there is some phenomenon affected the signal and limits the distance like attenuation, dispersion, and noise.

iii. Amplifiers, switches etc modify signal.

iv. Receiver decodes signal at end; here the most important factors is error rate or signal to noise ratio.

v. Bandwidth it defends as bits per second or megahertz/gigahertz; it’s increased by increasing data rate and number of optical in fiber; and limited by fiber dispersion, attenuation, noise, crosstalk, transmitter speed, receiver sensitivity and the range of wavelengths available.

2.2. **Optical Transmission System Concepts**

The basic components of an optical communication system shown in Figure 2.1, below can be summarized as follow [3]:

i. A serial bit stream in electrical form is presented to a modulator, which encodes the data appropriately for fiber transmission.

ii. A light source (laser or Light Emitting Diode - LED) is driven by the modulator and the light focused into the fiber.

iii. The light travels down the fiber (during which time it may experience dispersion and loss of strength).

iv. At the receiver end the light is fed to a detector and converted to electrical form.

v. The signal is then amplified and fed to another detector, which isolates the individual state changes and their timing. It then decodes the sequence of state changes and reconstructs the original bit stream (This overview is deliberately simplified. There are many ways to modulate the transmission and the details will vary from this example but the general principle remains unchanged).
vi. The timed bit stream so received may then be fed to a using device.

Figure 2.1: Optical Transmission – Schematic. [3]

2.3. Transmitting Light on a Fiber

An optical fiber is a very thin strand of silica glass in geometry. In reality it is a very narrow, very long glass cylinder with special characteristics. When light enters one end of the fiber it travels (confined within the fiber) until it leaves the fiber at the other end. Two critical factors stand out [3]:

i. Very little light is lost in its journey along the fiber.

ii. Fiber can bend around corners and the light will stay within it and be guided around the corners.

As shown in Figure 2.2, an optical fiber consists of two parts:

i. The core; which is a narrow cylindrical strand of glass.

ii. The cladding; which is tubular jacket surrounding the core.
The core has a (slightly) higher refractive index than the cladding. This means that the boundary (interface) between the core and the cladding acts as a perfect mirror. Light travelling along the core is confined by the mirror to stay within it - even when the fiber bends around a corner. When light is transmitted on a fiber, the most important consideration is “what kind of light?” The electromagnetic radiation that we call light exists at many wavelengths (Another way of saying this is that light has many frequencies or colors). These wavelengths go from invisible infrared through all the colors of the visible spectrum to invisible ultraviolet.

If a short pulse of light from a source such as a laser or an LED is sent down a narrow fiber, it will be changed (degraded) by its passage down the fiber. It will emerge (depending on the distance) much weaker, lengthened in time (“spearred out”), and distorted in other ways. There are some reason can causes this attenuation, Maximum Power, Polarization and Dispersion.

As this research concern in dispersion and how affected single mode fiber communications; so will going throw it for further clarification. The most important things related to this topic are [3- 3].
• Chirp Definition

Immediately after power is applied to a laser there is an abrupt change in the carrier (electron and hole) flux density in the cavity caused by the lasing operation itself. This density of charge carriers is one factor that affects the refractive index. In addition, the temperature in the cavity increases quite rapidly. This temperature increase is too localized to affect the length of the cavity immediately but it does contribute to changing the refractive index of the material in the active region (within the cavity). These changes in the RI of the cavity produce a rapid change in the centre wavelength of the signal produced. In the case of semiconductor lasers a “downward” chirp is produced. The wavelength shifts to a longer wavelength than it was immediately at the start of the pulse. It is not a large problem in short distances single-channel transmissions but in long distance applications and in WDM systems chirp can be a very serious problem. This is due to the fact that it broadens the spectral width of the signal and interacts with other aspects of the transmission system to produce distortion. Indeed, the chirp problem is the main reason that people use external modulators for transmission rates in excess of 1 Gbps [4].

• Wavelength Division Multiplexing

Wavelength Division Multiplexing (WDM) is the basic technology of optical networking. It is a technique for using a fiber (or optical device) to carry many separate and independent optical channels. The principle is identical to that used when we tune our television receiver to one of many TV channels. Each channel is transmitted at a different radio frequency and we select between them using a “tuner” which is just a resonant circuit within the TV set. Of course wavelength in the optical world is just the way we choose to refer to frequency and optical WDM is quite identical to radio FDM. Another way of envisaging WDM is to consider that each channel consists of light of a different color. Thus a WDM system transmits a “rainbow”. Actually at the wavelengths involved the light is invisible but it’s a good way of describing the principle [3].
There are many varieties of WDM. A simple form can be constructed using 1310 nm as one wavelength and 1550 as the other or 850 and 1310. This type of WDM can be built using relatively simple and inexpensive components and some applications have been in operation for a number of years using this principle. WDM is the basic technology for full optical networking [3].

2.4. Dispersion

- Dispersion Concept
  
i. Pulse spreading with distance
  ii. Depends on wavelength
  iii. Degrades bandwidth - pulses overlap.
  iv. The dispersion amount depends on type of fiber, possible of compensation and signal regeneration [4-3].

Figure 2.4: Effect of Dispersion [3]; (a) Short Distance, (b) Long Distance
The circles in the figure represent fiber loops. This is the conventional way to indicate distance in system diagrams.
• **Dispersion Mechanism**

Dispersion occurs when a pulse of light is spread out during transmission on the fiber. A short pulse becomes longer and ultimately joins with the pulse behind, making recovery of a reliable bit stream impossible. (In most communications systems bits of information are sent as pulses of light. 1 = light, 0 = dark. But even in analog transmission systems where information is sent as a continuous series of changes in the signal, dispersion causes distortion.) [4- 3].

• **Pulse Dispersion:**
  i. Depends on modal properties of fiber.
  ii. Determined by fiber type & composition.
  iii. Degree of dispersion depends on fiber type.

• **Effect of Dispersion**

Dispersion of transmitted optical signal causes distortion for both digital and analog transmission along optical fiber. When considering the major implementation of optical fiber transmission which involves some form of digital modulation, the dispersion mechanism within the fiber cause broadening of transmitted light pulses as they travel along the channel. It may be observed from the Figure (2.5) that each pulse broadens and overlaps with its neighbors, eventually becoming indistinguishable at the receiver input. The effect is known as inter symbol interference (ISI); thus an increasing number of errors may be encountered on the digital optical channel as the ISI becomes more pronounced [3];Figure 2.5 an illustration of the broadening of light pulses as they are transmitted along a fiber using the digital bit pattern 1011[3].
Figure 2.5: Broadening of Light Pulses; (a) Fiber Input, (b) Fiber Output at a Distance L1; (c) Fiber Output at a Distance L2>L1[3]
There are many kinds of dispersion, each of which works in a different way, but the most important three are discussed below [3]:

i. Material dispersion (chromatic dispersion); (ps/nm bandwidth, fiber/km); increases with source bandwidth; both lasers and LEDs produce a range of optical wavelengths (a band of light) rather than a single narrow wavelength. The fiber has different refractive index characteristics at different wavelengths and therefore each wavelength will travel at a different speed in the fiber. Thus, some wavelengths arrive before others and a signal pulse disperses (or spreads out).
ii. Modal dispersion (nanoseconds/fiber km); largest and depends on number of modes; when using multimode fiber (doesn’t occur at single mode fiber), the light is able to take many different paths or “modes” as it travels within the fiber. The distance traveled by light in each mode is different from the distance travelled in other modes. When a pulse is sent, parts of that pulse (rays or quanta) take many different modes (usually all available modes). Therefore, some components of the pulse will arrive before others. The difference between the arrival times of light taking the fastest mode versus the slowest obviously gets greater as the distance gets greater.

iii. Waveguide dispersion (ps/nm bandwidth, fiber km); is a very complex effect and is caused by the shape and index profile of the fiber core. However, this can be controlled by careful design and, in fact; waveguide dispersion can be used to counteract material dispersion as will be seen later.

2.5. Single-Mode Fiber

When the fiber core is very narrow compared to the wavelength of the light in use then the light cannot travel in different modes and thus the fiber is called “single-mode” or “monomode”. There is no longer any reflection from the core-cladding boundary but rather the electromagnetic wave is tightly held to travel down the axis of the fiber. It seems obvious that the longer the wavelength of light in use, the larger the diameter of fiber we can use and still have light travel in a single-mode. The core diameter used in a typical single-mode fiber is nine microns. It is not quite as simple as this in practice. A significant proportion (up to 20%) of the light in a single-mode fiber actually travels in the cladding. For this reason the “apparent diameter” of the core (the region in which most of the light travels) is somewhat wider than the core itself. The region in which light travels in a single-mode fiber is often called the “mode field” and the mode field diameter is quoted instead of the core diameter. The mode field varies in diameter depending on the relative refractive indices of core and cladding, Core diameter is a compromise. We can’t make the core too narrow because of losses at bends in the fiber. As the core diameter decreases compared to the wavelength (the core gets narrower or the wavelength gets longer), the minimum radius that we can bend the fiber without loss
increases. If a bend is too sharp, the light just comes out of the core into the outer parts of the cladding and is lost. The more important issue to make fiber single-mode is:

i. Making the core thin enough.
ii. Making the refractive index difference between core and cladding small enough.
iii. Using a longer wavelength.

Figure 2.7: Single-Mode Fiber [3]

This figure is not to scale. The core diameter is typically between 8 and 9 microns while the diameter of the cladding is 125 microns. Single-mode fiber usually has significantly lower attenuation than multimode (about half). Single-mode fibers have a significantly smaller difference in refractive index between core and cladding. The most important feature for the single mode fiber is [3]:

i. Single-mode transmission is simple.
ii. No modal dispersion.
iii. No modal noise.
iv. Transmission distance limited by chromatic dispersion.
v. Several types available differ in dispersion properties.

It’s not strictly correct to talk about “single-mode fiber” and “multimode fiber” without qualifying it - although we do this all the time. A fiber is single-moded or multimoded at a particular wavelength. If we use very long wave light (say 10.6 nm from a CO2 laser) then even most MM fiber would be single-moded for that wavelength. If we use 600 nm light on standard single-mode fiber then we do have a greater number of modes than just one (although typically only about 3 to 5). There is a single-mode fiber characteristic called the “cutoff wavelength”. This is typically around 1100 nm for single-mode fiber with a core diameter of 9 microns. The cutoff wavelength is the shortest wavelength at which the fiber remains single-moded. At wavelengths shorter than the cutoff the fiber is multimode. When light is introduced to the end of a fiber there
is a critical angle of acceptance. Light entering at a greater angle passes into the cladding and is lost. At a smaller angle the light travels down the fiber. If this is considered in three dimensions, a cone is formed around the end of the fiber within which all rays are contained. The sine of this angle is called the “numerical aperture” and is one of the important characteristics of a given fiber. Single-mode fiber has a core diameter of 4 to 10 μm (8 μm is typical). Multimode fiber can have many core diameters but in the last few years the core diameter of 62.5 μm in the US and 50 μm outside the US has become predominant.

- **Types of Single-Mode Fiber**

1. Step-Index single-mode fiber (simple)[3]:-
   i. Matched cladding
   ii. Cladding can be depressed by reduces core doping and dopes cladding to reduce index.

2. Dispersion-shifted fiber [3]:-
   i. More complex core-cladding design
   ii. Larger waveguide dispersion
   iii. Shifts zero chromatic dispersion
   iv. Higher dispersion slope.
   v. Minimum dispersion needed to prevent four-wave mixing.

The characteristics of single-mode fiber (Standard Fiber) were specified by the International Telecommunications Union (ITU) in the 1980’s. The key specifications are as follows [4-3]:

i. Cladding diameter = 125 microns.
ii. Mode field diameter = 9 10 microns at 1300 nm wavelength.
iii. Cutoff wavelength = 1100-1280 nm.
iv. Bend loss (at 1550 nm) must be less than 1 dB for travel through 100 turns of fiber wound on a spool of 7.5 cm diameter.
v. Dispersion in the 1300 nm band (1285-1330 nm) must be less than 3.5 ps/nm/km. At wavelengths around 1550 nm dispersion should be less than 20 ps/nm/km. (Picoseconds of dispersion per nanometer of signal bandwidth per kilometer of distance travelled).
vi. The rate of change of dispersion with wavelength must be less than .095 ps/nm²/km. This is called the dispersion slope.

2.6. Dispersion in Single-Mode Fiber

Since modal dispersion cannot occur in single-mode fiber (as only have one mode), the major sources of dispersion are chromatic dispersion (waveguide and material dispersion) and group velocity dispersion [3].

- Chromatic Dispersion

It’s a sum of material and waveguide dispersion. The two forms of dispersion have opposite signs, so they tend to counteract one another. Figure 2.8 shows the wavelength dependent dispersion characteristics of “standard” single-mode fiber. Notice that the two forms of dispersion cancel one another at a wavelength of 1310 nm. Thus if the signal is sent at 1310 nm dispersion will be minimized [3].

- Material Dispersion

This is caused by the fact that the refractive index of the glass we are using varies (slightly) with the wavelength. Some wavelengths therefore have higher group velocities and so travel faster than others. Since every pulse consists of a range of wavelengths it will spread out to some degree during its travel [3].

- Waveguide Dispersion

The shape (profile) of the fiber has a very significant effect on the group velocity. This is because the electric and magnetic fields that constitute the pulse of light extend outside of the core into the cladding. The amount that the fields overlap between core and cladding depends strongly on the wavelength. The longer the wavelength the further the electromagnetic wave extends into the cladding. The Refractive Index experienced by the wave is an average of the RI of core and cladding depending on the relative proportion of the wave that travels there. Thus since a greater proportion of the wave at shorter wavelengths is confined within the core, the shorter wavelengths “see” a higher RI than do longer wavelengths. (Because the RI of the core is higher than that of the
cladding). Therefore shorter wavelengths tend to travel more slowly than longer ones. Thus signals are dispersed (because every signal consists of a range of wavelengths) [3].

Figure 2.8: Step-Index Single-Mode Dispersion [9]

Figure 2.9: Dispersion-Shifted Single-Mode [9]
Dispersion (from all causes) is often grouped under the name GVD.

- **Group Velocity Dispersion**

  The group velocity of a wave is the velocity with which the overall shape of the wave's amplitudes — known as the modulation or envelope of the wave — propagates through space. For transmission system operation the most important & useful type of velocity is the group velocity, this is the actual velocity which the signal information & energy is traveling down the fiber. It is always less than the speed of light in the medium. Group velocity is the usual way of discussing the speed of propagation on a fiber. It is the speed of propagation of modulations along the fiber. It is generally a little less than the phase velocity. The reason that group velocity is different from phase velocity is related to the amount of dispersion of the medium. If there is no dispersion in the medium then group velocity and phase velocity are the same. The group velocity associated with the fundamental mode is frequency dependent because of chromatic dispersion. As a result, different spectral components of the pulse travel at slightly different group velocities, a phenomenon referred to as group-velocity dispersion (GVD), intermodal dispersion, or simply fiber dispersion. Intermodal dispersion has two contributions, material dispersion and waveguide dispersion. We consider both of them and discuss how GVD limits the performance of lightwave systems employing single-mode fibers. On standard single-mode fiber we consider two GVD regimes [3]:

  i. Normal Dispersion Regime is represented in Figure 2.9 at the left of the point where the line crosses the zero dispersion point. In this region the long wavelengths travel faster than the short ones! Thus after travelling on fiber wavelengths at the red end of the pulse spectrum will arrive first. This is called a positive chirp.

  **Note:** The use of the terms “red end” and “blue end” here requires some explanation. Any wavelength longer than about 700 nm is either visible red or infra-red. Thus all of the wavelengths in question can be considered “red”. However, it is very useful to identify the shorter wavelength (higher frequency) end of a pulse spectrum as the “blue end” and the longer wavelength (lower frequency) end as the “red end”.

  ii. Anomalous Dispersion Regime this is represented by the section of the figure to the right of the zero crossing point. Here, the short wavelengths (blue end of
the spectrum) travel faster than the long wavelengths (red end). After travel on a fiber the shorter wavelengths will arrive first. This is considered a negative chirp.

**Note:** The definitions of the words “normal” and “anomalous” given above are consistent with those used in most textbooks and in the professional literature. In some engineering contexts the use of the terms is reversed. That is, what we have defined above as normal becomes anomalous and what we defined as anomalous is called normal [3].

![Diagram of Fiber Dispersion (ps/nm/km)](image)

It seems obvious that the wider the spectral width of our signal the more dispersion we will have. Conversely, the narrower the signal the less dispersion. In SM fiber dispersion usually quoted in picoseconds of dispersion per nanometer of spectral width per kilometer of propagation distance (ps/nm/km) [3].
• Group Velocity & Group Delay [3]

The group velocity is given by [4]:

$$V_g = \frac{d\omega}{d\beta} \tag{2.1}$$

The group delay is given by:

$$\tau_g = \frac{\tau}{V_g} = \frac{t}{V_g} \frac{d\beta}{d\omega} \tag{2.2}$$

Where

$$V_g = \text{Group Velocity} ; \quad \omega = \text{Frequency} ; \quad \beta = \text{Propagation mode}$$

$$\tau_g = \text{Group delay} ; \quad \tau = \text{fiber optic length}$$

It is important to note that all above quantities depend both on frequency & the propagation mode.

2.7 Control of Dispersion in Single-Mode Fiber Links

Dispersion broadens a pulse by an amount unrelated to the length of the pulse. Dispersion becomes a problem for a receiver when it exceeds about 20% of the pulse length. Thus, if a pulse at 200 Mbps is dispersed on a given link by 15% then the system will probably work. If the data rate is doubled to 400 Mbps the dispersion will be 30% and the system will probably not work. Hence the higher data rate, the more important the control of dispersion becomes [3].

i. Modal dispersion does not exist in single-mode fiber. There is however, a trivial form of modal dispersion caused by birefringent effects spreading the two orthogonal polarization modes in a “single” mode fiber. This is called Polarization mode dispersion. However, the effect is usually trivial. In very short single-mode links (less than a few hundred meters) you can get modal dispersion due to additional modes being carried in the cladding. These disappear after a relatively short distance but they can be excited at the laser coupling or in a bad coupler or join.
ii. Material dispersion is significant in both types of fiber.

iii. Waveguide dispersion is significant in both Multimode and Single mode fibers but dominates in the SM case because there is no modal dispersion here.

Both material and waveguide dispersion are wavelength dependent effects. If we had a zero spectral width there would be no problem with these types of dispersion. Waveguide dispersion can be manipulated so that it acts in the opposite direction (has the opposite sign) to material dispersion. Single-mode fibers (for wide-area applications) of the late 1980s were adjusted such that the two forms of dispersion cancelled each other out at a wavelength of 1310 nm. For this reason, the 1300 nm band was widely used for long distance communication links at that time. However, the attenuation in the 1300 nm band is almost twice that of attenuation in the 1500 nm band. Worse, Erbium Doped Fiber Amplifiers (EDFAs) only work in the 1500 nm band, so if we want to use amplifiers, then we must use 1500 nm. Many things can be done to the fiber to reduce waveguide dispersion (such as varying the refractive index of core and cladding and changing the geometry of the fiber) and it is now possible to balance the two forms of dispersion at 1500 nm. This type of fiber is called Dispersion Shifted Fiber (DSF).

Another way of minimizing dispersion (both material and waveguide) is to use a narrow spectral width laser. These techniques combined have meant that almost all new long distance single-mode systems are being installed at 1500 nm [3-3].

**Dispersion Calculating**

Waveguide dispersion is usually quoted in (ps per nm per km) at a given wavelength. At 1500 nm a typical dispersion figure is 17 ps/nm/km. That is, a pulse (regardless of its length) will disperse by 17 picoseconds per nanometer of spectral width per kilometer of distance travelled in the fiber. So, in a typical single-mode fiber using a laser with a spectral width of 6 nm over a distance of 10 km we have [3]:

\[
\text{Dispersion} = 17\text{ps/nm/km} \times 6\text{nm} \times 10\text{km} = 1020\text{ps}
\]

At 1 Gbps a pulse is 1 ns long. So if we tried to send data over the above link at 1 Gbps then we would get 102% dispersion - that is, the system would not work. (20% is a good guideline for the acceptable limit.) But it would probably work quite well at a data
rate of 155 Mbps (a pulse length of 6.5 ns). A narrow spectral width laser might produce only one line with a line width of 300 MHz. Modulating it at 1 Gbps will add 2 GHz. 2,300 MHz is just less than .02 nm (at 1500 nm); that give [3]:

\[
\text{Dispersion} = 17 \text{ps/nm/km} \times 0.02 \text{nm} \times 10 \text{km} = 3.4 \text{ps}
\]

So in this point, dispersion just ceased to be a problem.

- **Control of Spectral Width**

Perhaps the most obvious thing we can do about dispersion is to control the spectral width of the signal! Chromatic dispersion is a linear function of spectral width. If you double the spectral width you double the dispersion. An important factor is that modulation adds to the bandwidth of the signal! The modulation broadens the signal by twice the highest frequency present in the modulating signal. Modulation with a square wave implies the presence of significant harmonics up to 5 times the fundamental frequency of the square wave! (Indeed a perfect square wave theoretically contains infinity of higher frequency components.); as example, if we want to modulate at 1 Gbps then the fundamental frequency is 500 MHz. A significant harmonic at 2.5 GHz will be present and therefore the broadening of the signal will be 5 GHz or about .04 nm. If we want to modulate at 10 GHz then signal broadening will be perhaps .4 nm. It is easily seen that these amounts are not significant if the laser spectral width is 5 nm but critically significant if the spectral width is .01 nm! This can be controlled by filtering the square wave modulating signal to remove higher frequency harmonics. But this filtering reduces the quality of the signal at the receiver. In practical systems we don’t worry about the 5th harmonic and usually can be content with the 3rd. So if we filter a 1 Gbps signal at about 1.5 GHz (at the transmitter) then we can usually build a receiver to suit [3].

- **Dispersion Shifted Fiber**

Is designed with dispersion zero point at around 1550 nm. For operation in the 1550 nm band this should be ideal. However, it is not always possible or indeed desirable.

i. In many cases we can’t have DSF because the fiber we must use is already installed. Digging up a few hundred kilometers of roadway to replace fiber types is an extremely costly exercise.
ii. If we are using (or planning to use) WDM technology the problems of four-wave mixing effectively prohibit the use of DSF.

iii. If we have a very long amplified link with many cascaded amplifiers we have another problem. The amplifiers will generate a certain amount of amplified spontaneous emission (ASE) noise at wavelengths near to the signal. While this ASE can be filtered out at the receiver it will usually be present on the link. Any ASE within about 2 nm in wavelength of the signal will undergo 4-wave mixing with the signal and create significant noise! Of course we could filter it out at the output stage of each amplifier but that would mean a long series of cascaded filters which would narrow the signal itself.

Except in the case of a limited number of amplifier spans DSF is not a good solution to the dispersion problem [4].

• Dispersion Compensating Fiber

For “Dispersion in Single-Mode Fiber” the core profile of a fiber can be controlled to produce just the amount of dispersion we want. In order to equalize an installed link with dispersion at 1550 nm of 17 ps/nm/km (standard fiber) we can connect a (shorter) length of compensating fiber in series with it. The compensating fiber typically has a dispersion of -100 ps/nm/km in the 1550 nm wavelength band. Because the dispersion acts in the opposite direction to the dispersion of the standard fiber the compensating fiber “undisperses” the signal. You might compensate a 100 km length of standard fiber for operation at 1550 nm by connecting it to 17 km of shifted fiber. However, almost by definition you are not installing the fiber new. So the added length of fiber sits at one end of the link on a drum. This adds to attenuation and additional amplification may be needed to compensate for the compensating fiber! DCF has a typical attenuation of .5 dB/km. In addition the narrow core of DCF makes it more susceptible to non-linear high power effects than standard fiber and it is also polarization sensitive[3].
The Faraday rotating mirror is used to rotate the polarization of the reflected signal. Thus any PMD is transit the same section of DCF again with rotated polarization. In addition, “undone” in the transit in the opposite direction. When installing a new optical fiber link there is a ability to plan for dispersion and the necessary compensation.

Figure 2.12 shows a link configuration with “lumped” dispersion compensation at the mid-span point of the link. In this case you only need half the length of DCF that you might need otherwise (because the light transits the DCF twice). The lumped DCF compensates for dispersion over the whole length of the link. This configuration has the advantage that you can use a Faraday rotating mirror to rotate the polarization of the signal. Thus any unwanted polarization dependencies introduced in the DCF are undone by the fact that the light has to transit the same section of DCF again with rotated polarization. In addition, unwanted polarization dependencies in the long link itself can be offset somewhat (to the extent that they are the same over the two halves of the link). Indeed this last point is the reason for sitting the DCF at mid-span. The major problem with this configuration is that you need access to the link at mid-span. This may not be easily possible in an installed link[3].
• **Balancing Dispersion on a Link**

Of course, if we are planning to operate in the 1550 nm band we could install Dispersion Shifted Fiber (DSF). This has a dispersion of zero at 1550 nm. However, as mentioned before, WDM systems have a severe problem if the fiber dispersion is really zero! This problem is called 4-wave mixing “Four-Wave Mixing (FWM)”. It turns out that for WDM operation needed some dispersion to minimize 4-wave mixing. Fiber with a dispersion of 4 ps/nm/km can be used to mitigate FWM but in very long amplified links (such as many undersea cables) even this minimal level of dispersion is a limitation. In this case the system architects sometimes employ a balanced structure where sections of dispersive fiber with different dispersion characteristics are joined to form the span. The idea here is that no section of fiber has zero dispersion but that different sections have dispersion of opposite sign so that the total at the end of the link (span) is zero.

In new link designers tend not to use such strongly dispersive fibers and instead might use a fiber with a dispersion of -2 ps/nm/km for the majority of the link. To compensate for this at intervals they insert a section (or sections) of standard (17 ps/nm/km) fiber in the link. There are a number of very long undersea links currently in operation which use this dispersion management technique. An undersea link with four WDM channels each operating at 2.4 Gbps over a distance of 4000 km has been reported as an operational system[3].

![Figure 2.13: Dispersion Compensation of a New Link with DCF [3]](image-url)
Note: In a WDM system it is quite hard to balance dispersive properties of fibers in this way. This is because the range over which the WDM signals are spread may be of the order of 30 or 40 nm. The dispersion characteristics of each fiber used will be different at different points over the wavelength range. Matching dispersion characteristics over a range of wavelengths can be very difficult. This may well result in one channel of the WDM spectrum having zero dispersion (total at the end of the link) and other channels having significant finite dispersion![3-3].

- **Mid-Span Spectral Inversion**

The concept here is to use a device in the middle of the link to invert the spectrum. This process changes the short wavelengths to long ones and the long wavelengths to short ones. If you invert the spectrum in the middle of a link (using standard fiber) the second half of the link acts in the opposite direction (really the same direction but the input has been exactly pre-emphasized). When the pulse arrives it has been re-built exactly - compensated for by the second half of the fiber. Mid-span spectral inversion is a bit difficult to implement in all situations because you have to put an active device into the middle of the fiber link. This may or may not be practical (you might not be able to get access to the mid-point of the fiber link) [3].

![Spectral Inverter – Schematic Dispersion](image)

Figure 2.14: Spectral Inverter – Schematic Dispersion [3]

Figure 2.15 shown a dispersed pulse before input to a phase conjugation process and then at output. The wavelength spectrum has been completely inverted. This spectral inversion is performed by a process called “optical phase conjugation”. Devices that change the wavelength using either 4-Wave Mixing or Difference Frequency Generation invert the spectrum as a byproduct of their wavelength conversion function. These can be used as spectral inverters if we can tolerate the wavelength shift involved. Although there are devices that can perform phase conjugation (the spectral inversion function) in reality what is often used here is just a section of dispersion compensating fiber on a
drum. Several kilometers are typically used. So this can reduce to just another configuration option of dispersion compensating fiber [3].

![Figure 2.15: Spectral Inversion - Wavelength (Frequency) Domain [3]](image)

- **Chirped Fiber Bragg Gratings**

  Fiber Bragg Gratings (FBGs) are perhaps the most promising technology for dispersion compensation. A “chirped” FBG is used where the spacing’s of the lines on the grating vary continuously over a small range. Shorter wavelength light entering the grating travels along it almost to the end before being reflected. Longer wavelength light is reflected close to the start of the grating. Thus short wavelengths are delayed in relation to longer ones. Since the pulse has been dispersed such that short wavelengths arrive before the long ones, the grating can restore the original pulse shape. It undoes the effects of dispersion. The configuration requires a circulator to direct the light in and out of the grating as shown in Figure 2.16. Chirped FBGs need to be quite long (for a simple single-channel application up to 20 cm is commonly required). In a WDM system a fully continuous chirp would require a very long grating indeed. To compensate for 100 km of standard (17 ps/nm/km) fiber the chirped grating needs to be 17 cm long for every nm of signal bandwidth! In this instance a WDM system with channels spread over (say) 20 nm would need a chirped FBG (20 × 17) 340 cm long! Long FBGs are very hard to construct! The current technological limit is about 1 meter in length. In a single-channel application you build the grating as a concatenated series of fixed wavelength gratings rather than a continuous variable grating. This means you don’t need to maintain phase continuity between the different sections. You can write the grating in a number of separate operations and of course this means that it can be as long as you like[3].
The major problem with chirped FBGs is that they have a ripple characteristic in the GVD they produce. (The aim of a chirped FBG is to produce GVD in the opposite direction from that produced on the fiber link.) This ripple can be a source of transmission system noise. The longer the grating the larger the problem with ripple and its resultant noise. In addition short FBGs are filters. When you process a signal through many stages of filtering you get a very narrow signal as a result and this can also distort and add noise to the signal.

2.7 Dispersion Management

It should be clear that with the advent of optical amplifiers, fiber losses are no longer a major limiting factor for optical communication systems. Indeed, modern lightwave systems are often limited by the dispersive and nonlinear effects rather than fiber losses. In some sense, optical amplifiers solve the loss problem but, at the same time, worsen the dispersion problem since, in contrast with electronic regenerators; an optical amplifier does not restore the amplified signal to its original state.

As a result, dispersion-induced degradation of the transmitted signal accumulates over multiple amplifiers. For this reason, several dispersion-management schemes were developed during the 1990s to address the dispersion problem. The limitations imposed on the system performance by dispersion-induced pulse broadening, the group-velocity dispersion (GVD) effects can be minimized using a narrow-linewidth laser and operating close to the zero-dispersion wavelength $\lambda_{ZD}$ of the fiber. However, it is not always practical to match the operating wavelength $\lambda$ with $\lambda_{ZD}$. Such systems generally use the
existing fiber-cable network installed during the 1980s and consisting of more than 50 million kilometers of the “standard” single-mode fiber with $\lambda_{ZD} = 1.31 \mu m$ since the dispersion parameter $D \approx 16 \text{ps/(km} - \text{nm)}$ in the $1.55 - \mu m$ region of such fibers, the GVD severely limits the performance when the bit rate exceeds 2 GB/s[4].

![Figure 2.17: Limiting Bit Rate of Single-Mode Fibers[4]](image)

Figure 2.17 shown the limiting bit rate of single-mode fibers as a function of the fiber length for $\sigma_{\lambda}=0,1, \text{and } 5 \text{ nm}$ the case $\sigma_{\lambda}=0$ corresponds to the case of an optical source whose spectral width is much smaller than the bit rate [4].

Even when $L$ increases by a factor of 10 because of the $L^{-1/3}$ dependence of the bit rate on the fiber length. The dashed line in Fig. 2.19 shows this dependence by using Eq. $B(|\beta_3|L)^{1/3} \leq 0.324$ with $\beta_3 = 0.1 \text{ps}^3/\text{km}$. Clearly, the performance of a lightwave system can be improved considerably by operating it near the zero-dispersion wavelength of the fiber and using optical sources with a relatively narrow spectral width. For a directly modulated DFB laser, we can use Eq. $BL|D|\sigma_3 \leq \frac{1}{4}$ for estimating the maximum transmission distance so that the [4].
\[ L < (4B|D|S_\lambda)^{-1} \] (2.3)

Where \( S_\lambda \) is the root-mean-square (RMS) width of the pulse spectrum broadened considerably by frequency chirping. Using \( D = 16 \text{ ps/(km-nm)} \) and \( S_\lambda = 0.15 \text{ nm} \) in Eq(2.3 ), lightwave systems operating at 2.5 Gb/s are limited to \( L \approx 42 \text{ km} \). Indeed, such systems use electronic regenerators, spaced apart by about 40 km, and cannot benefit from the availability of optical amplifiers. Furthermore, their bit rate cannot be increased beyond 2.5 Gb/s because the regenerator spacing becomes too small to be feasible economically [4].

System performance can be improved considerably by using an external modulator and thus avoiding spectral broadening induced by frequency chirping. This option has become practical with the commercialization of transmitters containing DFB lasers with a monolithically integrated modulator. The \( S_\lambda = 0 \) line in Fig. 2.19 provides the dispersion limit when such transmitters are used with the standard fibers. The limiting transmission distance is then obtained by using the following Eq [4]:

\[ \sigma^2 = \sigma_0^2 + (\beta_3L/4\sigma_0^2)^2/2 \equiv \sigma_0^2 + \sigma_Z^2 \]

And is given by [4]:

\[ L < (16|\beta_2|B^2)^{-1} \]

(2.4)

Where \( \beta_2 \) is the GVD coefficient related to \( D \) by [4]:

\[ D = \frac{d}{d\lambda} \left( \frac{1}{v_g} \right) = -\frac{2\pi c}{\lambda^2} \beta_2. \]

Where

\( D = \text{Dispersion parameter} ; \quad \beta_2 = \text{GVD coefficient} ; \quad \lambda = \text{Wavelength} \)

If we use a typical value of \( \beta_2 = -20 \text{ ps}^2/\text{km} \) at 1.55 \( \mu \text{m} \), \( L < 500 \text{ km at 2.5 Gb/s} \) although improved considerably compared with the case of directly modulated DFB lasers, this dispersion limit becomes of concern when in-line amplifiers are used for loss compensation. Moreover, if the bit rate is increased to 10 Gb/s, the GVD-limited transmission distance drops to 30 km, a value so low that optical amplifiers cannot be used in designing such lightwave systems. It is evident from Eq. (2.4) that the relatively large GVD of standard single-mode fibers severely limits the performance of 1.55-\( \mu \text{m} \)
systems designed to use the existing telecommunication network at a bit rate of 10 Gb/s or more.

A dispersion-management scheme attempts to solve this practical problem. The basic idea behind all such schemes is quite simple and can be understood by using the pulse-propagation equation [4] and written as:

\[
\frac{\partial A}{\partial z} + \frac{i\beta_3}{\alpha^2} \frac{\partial^2 A}{\partial t^2} + \frac{\beta_4}{6} \frac{\partial^3 A}{\partial t^3} = 0
\] (2.5)

Where A is the pulse-envelope amplitude. The effects of third-order dispersion are included by the \(\frac{\partial^3 A}{\partial t^3}\) term. In practice, this term can be neglected when \(|\frac{\partial^2 A}{\partial t^2}| > 0.1 \frac{G_{pp}}{mm} \) Eq (2.5) has been solved [4] and the solution is given by [4]:

\[
A(z, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{A}(0, \omega) \exp \left( \frac{i}{2} \beta_2 \omega^2 + \frac{i}{6} \beta_3 \omega^3 - i\omega t \right) d\omega,
\]

In the specific case of \(\beta_3 = 0\) the solution becomes:

\[
A(z, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{A}(0, \omega) \exp \left( \frac{i}{2} \beta_2 \omega^2 - i\omega t \right) d\omega,
\] (2.6)

Where \(\tilde{A}(0, \omega)\) is the Fourier transform of \(A(0, t)\).

Dispersion-induced degradation of the optical signal is caused by the phase factor \(\exp \left( \frac{i}{2} \beta_2 \omega^2 / 2 \right)\) acquired by spectral components of the pulse during its propagation in the fiber. All dispersion-management schemes attempt to cancel this phase factor so that the input signal can be restored. Actual implementation can be carried out at the transmitter, at the receiver, or along the fiber link [4].

- **Precompensation Schemes**

This approach to dispersion management modifies the characteristics of input pulses at the transmitter before they are launched into the fiber link. The underlying idea can be understood from Eq. (2.6). It consists of changing the spectral amplitude \(\tilde{A}(0, \omega)\) of the input pulse in such a way that GVD-induced degradation is eliminated, or at least reduced substantially. Clearly, if the spectral amplitude is changed as [4]:

32
\[ \tilde{A}(0, \omega) \rightarrow \tilde{A}(0, \omega)exp(-i\omega^2 \beta_2 L/2) \] (2.7)

Where L is the fiber length, GVD will be compensated exactly, and the pulse will retain its shape at the fiber output. Unfortunately, it is not easy to implement Eq. (2.7) in practice. In a simple approach, the input pulse is chirped suitably to minimize the GVD-induced pulse broadening. Since the frequency chirp is applied at the transmitter before propagation of the pulse, this scheme is called the prechirp technique. The more important precompensation schemes are [4]:

i. Prechirp Technique.
ii. Novel Coding Techniques.
iii. Nonlinear Prechirp Techniques.

• Postcompensation Techniques

Electronic techniques can be used for compensation of GVD within the receiver. The philosophy behind this approach is that even though the optical signal has been degraded by GVD, one may be able to equalize the effects of dispersion electronically if the fiber acts as a linear system. It is relatively easy to compensate for dispersion if a heterodyne receiver is used for signal detection [4]. A heterodyne receiver first converts the optical signal into a microwave signal at the intermediate frequency \(\omega_{IF}\) while preserving both the amplitude and phase information. A microwave bandpass filter whose impulse response is governed by the transfer function [4]

\[ H(\omega) \exp[-i(\omega - \omega_{IF})^2 \beta_2 L/2] \] (2.8)
Figure 2.18: Dispersion-Limited Transmission Distance [4]

Figure 2.18 shown the dispersion-limited transmission distance as a function of launch power for Gaussian (m = 1) and super-Gaussian (m = 3) pulses at bit rates of 4 and 8 Gb/s. Horizontal lines correspond to the linear case [4]. Where L is the fiber length, should restore to its original form the signal received. This conclusion follows from the standard theory of linear systems [4] by using Eq. (2.6) with z=L. This technique is most practical for dispersion compensation in coherent lightwave systems. In a 1992 transmission experiment, a 31.5-cm-long microstrip line was used for dispersion equalization. Its use made it possible to transmit the 8-Gb/s signal over 188 km of standard fiber having a dispersion of 18.5 ps/(km-nm). In a 1993 experiment, the technique was extended to homodyne detection using single-sideband transmission, and the 6-Gb/s signal could be recovered at the receiver after propagating over 270 km of standard fiber. Microstrip lines can be designed to compensate for GVD acquired over fiber lengths as long as 4900 km for a lightwave system operating at a bit rate of 2.5 Gb/s[4].

Use of a coherent receiver is often not practical. An electronic dispersion equalizer is much more practical for a direct-detection receiver. A linear electronic circuit cannot compensate GVD in this case. The problem lies in the fact that all phase information is lost during direct detection as a photodetector responds to optical intensity only [4-4]. As
a result, no linear equalization technique can recover a signal that has spread outside its allocated bit slot. Nevertheless, several nonlinear equalization techniques have been developed that permit recovery of the degraded signal. In one method, the decision threshold, normally kept fixed at the center of the eye diagram [4-4], is varied depending on the preceding bits. In another, the decision about a given bit is made after examining the analog waveform over a multiple-bit interval surrounding the bit in question. The main difficulty with all such techniques is that they require electronic logic circuits, which must operate at the bit rate and whose complexity increases exponentially with the number of bits over which an optical pulse has spread because of GVD-induced pulse broadening. Consequently, electronic equalization is generally limited to low bit rates and to transmission distances of only a few dispersion lengths. An optoelectronic equalization technique based on a transversal filter has also been proposed. In this technique, a power splitter at the receiver splits the received optical signal into several branches. Fiber-optic delay lines introduce variable delays in different branches. The optical signal in each branch is converted into photocurrent by using variable-sensitivity photodetectors, and the summed photocurrent is used by the decision circuit. The technique can extend the transmission distance by about a factor of 3 for a lightwave system operating at 5 Gb/s[4-4].

- **Dispersion-Compensating Fibers**

  The preceding techniques may extend the transmission distance of a dispersion-limited system by a factor of 2 or so but are unsuitable for long-haul systems for which GVD must be compensated along the transmission line in a periodic fashion. What one needs for such systems is an all-optical, fiber-based, dispersion-management technique. A special kind of fiber, known as the dispersion-compensating fiber (DCF), has been developed for this purpose as mentioned before. The use of DCF provides an all-optical technique that is capable of compensating the fiber GVD completely if the average optical power is kept low enough that the nonlinear effects inside optical fibers are negligible. It takes advantage of the linear nature of Eq. (2.5). To understand the physics behind this dispersion-management technique, consider the situation in which each optical pulse propagates through two fiber segments, the second of which is the DCF. Using Eq. (2.6) for each fiber section consecutively, we obtain [4-4]:

35
\[ A(L, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \exp \left[ i \frac{\omega^2}{4} (\beta_{21} L_1 + \beta_{22} L_2) - i \omega t \right] d\omega \]  

(2.9)

Where:

\[ L = L_1 + L_2, \quad \text{and} \]

\[ \beta_{2j} \equiv \text{GVD parameter for the fiber segment of length } L_j \quad (j=1, 2). \]

If the DCF is chosen such that the \( \omega^2 \) phase term vanishes, the pulse will recover its original shape at the end of DCF. The condition for perfect dispersion compensation is thus \( \beta_{21} L_1 + \beta_{22} L_2 = 0 \), or [4]:

\[ D_1 L_1 + D_2 L_2 = 0 \]

(2.10)

Eq (2.10) shows that the DCF must have normal GVD at 1.55 \( \mu \)m \( D_2 < 0 \) because \( D_1 > 0 \) for standard telecommunication fibers. Moreover, its length should be chosen to satisfy [4]:

\[ L_2 = -(D_1 / D_2) L_1 \]

(2.11)

For practical reasons, \( L_2 \) should be as small as possible. This is possible only if the DCF has a large negative value of \( D_2 \). A practical solution for upgrading the terrestrial lightwave systems making use of the existing standard fibers consists of adding a DCF module (with 6–8 km of DCF) to optical amplifiers spaced apart by 60–80 km. The DCF compensates GVD while the amplifier takes care of fiber losses. This scheme is quite attractive but suffers from two problems. First, insertion losses of a DCF module typically exceed 5 dB. Insertion losses can be compensated by increasing the amplifier gain but only at the expense of enhanced ASE noise. Second, because of a relatively small mode diameter of DCFs, the effective mode area is only \( \sim 20 \mu m^2 \). As the optical intensity is larger inside a DCF at a given input power, the nonlinear effects are considerably enhanced. The problems associated with a DCF can be solved to a large extent by using a two mode fiber designed with values of \( V \) such that the higher-order mode is near cutoff (\( V \approx 2.5 \)). Such fibers have almost the same loss as the single-mode fiber but can be designed such that the dispersion parameter \( D \) for the higher-order mode has large negative values. Indeed, values of \( D \) as large as \(-770 \) ps/(km-nm) have been measured for elliptical-core fibers. A 1-km length of such a DCF can compensate the
GVD for a 40-km-long fiber link, adding relatively little to the total link loss. The use of a two-mode DCF requires a mode-conversion device capable of converting the energy from the fundamental mode to the higher-order mode supported by the DCF. Several such all-fiber devices have been developed. The all-fiber nature of the mode-conversion device is important from the standpoint of compatibility with the fiber network. Moreover, such an approach reduces the insertion loss. Additional requirements on a mode converter are that it should be polarization insensitive and should operate over a broad bandwidth. Almost all practical mode-conversion devices use a two-mode fiber with a fiber grating that provides coupling between the two modes. The grating period \( \Lambda \) is chosen to match the mode-index difference \( \delta \bar{n} \) of the two modes (\( \Lambda = \lambda / \delta \bar{n} \)) and is typically 100 \( \mu \text{m} \). Such gratings are called long-period fiber gratings. Figure 2.21 shows schematically a two-mode DCF with two long-period gratings. The measured dispersion characteristics of this DCF are also shown. The parameter D has a value of \(-420 \text{ ps/(km-nm)}\) at 1550 nm and changes considerably with wavelength. This is an important feature that allows for broadband dispersion compensation. In general, DCFs are designed such that \(|D|\) increases with wavelength. The wavelength dependence of D plays an important role for wavelength-division multiplexed (WDM) systems [4].

![Figure 2.19](image.png)

(a) Schematic of a DCF made using a higher-order mode (HOM)

(b) Dispersion Spectrum of the DCF fiber and two long period’s gratings (LPGs) and Dispersion spectrum of the DCF.
Optical Filters

A shortcoming of DCFs is that a relatively long length (> 5 km) is required to compensate the GVD acquired over 50 km of standard fiber. This adds considerably to the link loss, especially in the case of long-haul applications. For this reason, several other all-optical schemes have been developed for dispersion management. Most of them can be classified under the category of optical equalizing filters. Interferometric filters are considered in this section. The function of optical filters is easily understood from Eq. (2.6). Since the GVD affects the optical signal through the spectral phase \( \exp(i\beta_2 z \omega^2)/2 \), it is evident that an optical filter whose transfer function cancels this phase will restore the signal. Unfortunately, no optical filter (except for an optical fiber) has a transfer function suitable for compensating the GVD exactly. Nevertheless, several optical filters have provided partial GVD compensation by mimicking the ideal transfer function. Consider an optical filter with the transfer function \( H(\omega) \). If this filter is placed after a fiber of length \( L \), the filtered optical signal can be written using Eq. (2.6) as [4]:

\[
A(L, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{A}(0, \omega)H(\omega)\exp\left(i\frac{4}{\beta_2}L\omega^2 - i\omega t\right) d\omega, \tag{2.12}
\]

Figure 2.20 shown dispersion management in a long-haul fiber link using optical filters after each amplifier. Filters compensate for GVD and also reduce amplifier noise [4].
• **Fiber Bragg Gratings**

A fiber Bragg grating acts as an optical filter because of the existence of a stop band, the frequency region in which most of the incident light is reflected back. The stop band is centered at the Bragg wavelength \( \lambda_B = 2\bar{n}\Lambda \), where \( \Lambda \) is the grating period and \( \bar{n} \) is the average mode index. The periodic nature of index variations couples the forward- and backward-propagating waves at wavelengths close to the Bragg wavelength and, as a result, provides frequency-dependent reflectivity to the incident signal over a bandwidth determined by the grating strength. In essence, a fiber grating acts as a reflection filter. Although the use of such gratings for dispersion compensation was proposed in the 1980s, it was only during the 1990s that fabrication technology advanced enough to make their use practical [4].

i. Uniform-Period Gratings.

ii. Chirped Fiber Gratings.

iii. Chirped Mode Couplers.

• **Optical Phase Conjugation**

Although the use of optical phase conjugation (OPC) for dispersion compensation was proposed in 1979, it was only in 1993 that the OPC technique was implemented experimentally; it has attracted considerable attention since then [4].

i. Principle of Operation.

ii. Compensation of Self-Phase Modulation.

iii. Phase-Conjugated Signal.

• **Long-Haul Lightwave Systems**

This section focused on lightwave systems in which dispersion management helps to extend the transmission distance from a value of \( \sim 10 \text{ km} \) to a few hundred kilometers. The important question is how dispersion management can be used for long-haul systems for which transmission distance is several thousand kilometers. If the optical signal is regenerated electronically every 100–200 km, all techniques discussed in this section should work well since the nonlinear effects do not accumulate over long lengths. In contrast, if the signal is maintained in the optical domain over the entire link by using periodic amplification, the nonlinear effects such as SPM, crossphase modulation
(XPM), and FWM would limit the system ultimately. Indeed, the impact of nonlinear effects on the performance of dispersion-managed systems has been a subject of intense study [4]:

i. Periodic Dispersion Maps.
ii. Simple Theory.
iii. Intrachannel Nonlinear Effects.

• High-Capacity Systems

Modern WDM lightwave systems use a large number of channels to realize a system capacity of more than 1 Tb/s. For such systems, the dispersion-management technique should be compatible with the broad bandwidth occupied by the multichannel signal [4].

i. Broadband Dispersion Compensation.
ii. Tunable Dispersion Compensation.
iii. Higher-Order Dispersion Management.
iv. PMD Compensation.
Chapter Three
Design and Implementation
3.1 Effects of Group Velocity Dispersion (GVD) on Gaussian Pulse Propagation.

To demonstrate the influence of the group (GVD) velocity dispersion on pulse propagation in optical fibers in "linear" regime. The basic effects related to GVD are[10]:

i. GVD induced pulse broadening.
ii. GVD induced pulse chirping.
iii. Pulse compression.

The equation, which describes the effect of GVD on optical pulse propagation neglecting the losses and nonlinearities, is [1]:

$$\frac{i}{t} \frac{\partial E}{\partial z} = \beta_2 \frac{\partial^2 E}{\partial t^2}$$  \hspace{1cm} (3.1)

Where is:

- $z$ = the propagation direction,
- $t$ = is the time,
- $E$ = is the electric field envelope, and
- $\beta_2 = \frac{\partial^2 \beta}{\partial \omega^2}$ is the GVD parameter, defined as the second derivative of the fiber mode propagation constant with value to frequency.

For an input pulse with a Gaussian shape,

$$E(z = 0, t) = \sqrt{p_0} \exp \left[ -\frac{t^2}{2T_0^2} \right]$$  \hspace{1cm} (3.2)

The pulse width $T_0$ (related to the pulse full width at half maximum by $T_{FWHM} \approx 1.665T_0$) increases with $z$ (the pulse broadens) according to[10]:

$$T(z) = \left[ 1 + \left( \frac{z}{L_D} \right)^2 \right]^{1/2} T_0$$  \hspace{1cm} (3.3)

And, consequently, the peak power changes, due to GVD, are given by:
p(\sigma) = \frac{p_0}{\left[1 + \left(\frac{\sigma}{L_D}\right)^2\right]^{1/2}} \quad (3.4)

In Equation (3.3) and Equation (3.4), the quantity $L_D = \frac{\tau_0^2}{|\beta_2|}$ is the dispersion length. Its meaning is quite straightforward: after propagating a distance equal to $L_D$, the pulse broadens by a factor of $\sqrt{2}$. To demonstrate this, we created the following simple circuit [1&10].

Figure 3.1: GVD Circuit Design [10]
The most important factor to set is:

i. Bit rate equal to 40 GB/S.

ii. Correspond to bit duration of 25 ps.

iii. Using default value of 0.5 for “width” of the Optical Gaussian Pulse Generator.

iv. The resulting FWHM of the pulse is 12.5 ps.

v. The $T_0$ parameter is then;

$$T_0 \approx \frac{T_{FWHM}}{1.665} = \frac{12.5\text{ps}}{1.665} = 7.5\text{ps}$$

And using the value of $\beta_2 \approx -20 (\text{ps})^2/\text{km}$ at 1.55\mu m for SMF the dispersion length is[10]:

$$L_D = \frac{T_0^2}{|\beta_2|} = \frac{(7.5^2)}{20} = 2.812\text{km}$$

In the Optical Fiber properties, we set the length of the fiber equal to this value, and we disable all the effects except GVD [10].
Figure 3.3: Optical Fiber Properties [10]
3.2 Compensation of Dispersion with Ideal Dispersion Component

The most major features for the Ideal Dispersion Component are [10]:

i. Large negative dispersion coefficient.
ii. Low attenuation.
iii. Minimal nonlinear contributions.
iv. Wide bandwidth.
v. Corrects dispersion slope as well.
vi. Minimal ripple.
vii. Polarization independent.
viii. Manufacturable.

To demonstrate the Compensation of dispersion with Ideal dispersion component in OpiSystem same components as described at last circuit with addition components as follow and shown in figure 3.4[10]:

i. Ideal dispersion component.
ii. Optical Time Domain Visualizer.
iii. Optical Spectrum Analyzer.

The most major important factor to set is[1&10]:

i. Bit rate = 40 Gb/s
ii. Bit duration = 12.5 ps pulse (initial of the Optical Gaussian pulse generator).
iii. Value for width = 0.5.
iv. Single mode fiber length = 10 km
v. Dispersion parameter = - 160 ps/nm
Figure 3.4 (a) Circuit Design, Compensation with Ideal Dispersion Component [10]

Figure 3.4 (b) Bill of Material for Dispersion, Compensation with Ideal Dispersion Component [10]
Chapter Four
Results and Discussion
4.1 Result of Demonstration the Effects of Group Velocity Dispersion (GVD) on Gaussian Pulse Propagation.

The project calculated and the obtained results are presented in Figure (4.1). The pulse had broadened (the peak power decreases in accordance with Equation (3.4)). The origin of pulse broadening can be understood by looking at the instant frequency of the pulse, namely the chirp [2].

Figure 4.1: Input Pulse and Output Pulse of Effect GVD on Gaussian Pulse Propagation.
(a) Optical Time Domain Visualizer_2 (Input).
Figure 4.1: Input Pulse and Output Pulse of Effect GVD on Gaussian Pulse Propagation.
(b) Bill of Material for Dispersion
Figure 4.2: Pulse Chirp Plotted with Intensity Pulse.
(a) Optical Time Domain Visualizer_2 (Input).
(b) Optical Time Domain Visualizer (Output) [2]

The straight lines represent the chirp[2].
As shown in Figure (4.2), where the pulse chirp is plotted together with the pulse intensity. Whereas the input pulse is chirpless, the instantaneous frequency of the output pulse decreases from the leading to the trailing edge of the pulse. The reason for this is GVD. In the case of anomalous GVD ($\beta_2 < 0$), the higher frequency ("blue-shifted") components of the pulse travel faster than the lower frequency (or "red-shifted") ones.

Note: The leading edge of the pulse is blue shifted and the trailing edge of the pulse is red-shifted. Because the "blue" and "red" spectral components tend to separate in time, this leads to pulse broadening. However, the pulse spectrum remains unchanged, as Figure (4.3) shown [2].

Figure 4.3: Spectra Corresponding to Figure 4.1 and 4.2
(a) Optical Spectrum Analyzer 1 (Input).
If the input pulse is frequency modulated (i.e. chirped), Equation (3.2) is replaced by:

\[
E(z = 0, t) = \sqrt{p_0} \exp \left(1 + \frac{1+iC}{2} \frac{t^2}{\tau_0^2}\right) \tag{4.1}
\]

And the expression for the dependence of the pulse width on \(z\) is [1]:

\[
T(z) = T_0 \left[1 + \frac{c \beta_2 z}{i \delta} \right]^2 + \left(\frac{\beta_2 z}{\tau_0^2} \right)^2 \right]^{1/2} \tag{4.2}
\]

The pulse broadens monotonically with \(z\) if \(\beta_2 C > 0\), however, it goes through initial narrowing when \(\beta_2 C < 0\) [1]:

\[
z_{\text{min}} = \frac{|C|}{1+C^2} L_D \tag{4.3}
\]

And is given by:

\[
T(Z_{\text{min}}) = \frac{T_0}{(1+C^2)^{1/2}} \tag{4.4}
\]
In this case the peak power of the pulse is [2]:

\[ P(z_{\text{min}}) = P_0 (1 + C^2)^{1/2} \]  \hspace{1cm} (4.5)

Initial narrowing of the pulse for the case \( \beta_2 C < 0 \) can be explained by noticing that in this case the frequency modulation (or "chirp") is such that the faster ("blue" in the case of anomalous GVD) frequency components are in the trailing edge, and the slower (or "red" in the case of anomalous GVD) in the leading edge of the pulse. As the pulse propagates, the faster components will overtake the slower ones, leading to pulse narrowing. At the same time, the dispersion induced chirp will compensate for the initial one. At \( z = z_{\text{min}} \), full compensation between both will occur. With further propagation, the fast and the slow frequency components will tend to separate in time from each other and, consequently, pulse broadening will be observed [2].

To demonstrate this, we use a chirped Gaussian pulse with the chirp parameter \( C = 2 \) (since \( \beta_2 < 0 \) in our case) Figure (4.4)[2].

![Optical Gaussian Pulse Generator Properties](OpticalGaussianPulseGeneratorProperties.png)

Figure 4.4: Setting the Chirp Parameter to Observe Pulse Compression [2]
Using Eq (4.3) and Eq (4.5); can obtained that:

\[ z_{\text{min}} = \frac{|C|}{1 + C^2} L_D = \frac{2}{5} L_D = 1.125 \text{km}; \]

\[ P(z_{\text{min}}) = P_0 (1 + C^2)^{1/2} = \sqrt{5} \approx 2.23 \text{ mW} \]

The length of the fiber changed to 1.125 km and calculate the project (the dispersion length is reduced to 1.125 km) [2].

![Figure 4.5: Pulse shape and chirp at z=0 and z=z_{\text{min}}. (a) Optical Time Domain Visualizer_2 (Input).](image)

Figure 4.5: Pulse shape and chirp at z=0 and z=z_{\text{min}}. (a) Optical Time Domain Visualizer_2 (Input).
Figure 4.5: Pulse shape and chirp at $z=0$ and $z=z_{\text{min}}$. (b) Optical Time Domain Visualizer (Output). (c) Optical Time Domain Visualizer Chirp (Input).
The results for the output pulse shape and chirp are presented in Figure (4.5). It can be seen that an exact compensation between the dispersion induced and initial chirp occurs, and that the peak power of the pulse is $\approx 2.23 \text{mW}$, as given by Eq (4.5) [1&2].

- Note: There is exact compensation between the initial and the dispersion-induced chirp [2].
4.2 Result of Demonstration Compensation of dispersion with Ideal Dispersion Component:

Figure 4.6: Initial Gaussian pulse and Gaussian pulse after 10 km propagation in SMF [2]
(a) Optical Time DomainVisualizer 1 (Input).
(b) Optical Time DomainVisualizer 2 (Output)
Figure (4.6) shown the results when the pulse was launched in 10 km SMF. As a result of this propagation, the width of the pulse increases approximately four times. Then if the main tab of the Ideal Dispersion Compensation component adjusted at (- 160 ps/nm). the output of the Ideal Dispersion Compensation as shown in Figure (4.7) [2].

![Figure 4.7: The Result of Dispersion Compensation Performed with the Ideal Dispersion Compensation Component [2]](image)

As expected, an exact compensation of accumulated dispersion was achieved. In conclusion, we have shown in this session how to use an Ideal Dispersion Compensation component in OptiSystem for dispersion compensation.

When some parameter is changed as mentioned below and takes other reading the result as show at figure as (4.8) [2]:

i. Single mode fiber length = 100 km.

ii. Dispersion parameter = - 1600 ps/nm.
Figure 4.8: Initial Gaussian Pulse and Gaussian Pulse after 100 km Propagation in SMF [2].

(a) Optical Time DomainVisualizer 1 (Input).
(b) Optical Time DomainVisualizer 2 (Output).
Figure 4.8: Initial Gaussian Pulse and Gaussian Pulse after 100 km Propagation in SMF [2].
(c) Optical Time Domain Visualizer 3 (Output)
Chapter Five
Conclusion and Recommendation.
5.1 Conclusion

- In this project, the properties of optical fibers were discussed. When the GVD was considered, and shown how it affected the signal and reduced the system performance and limiting the transmission distant.

- As mentioned dispersion represents a broad class of phenomena related to the fact that the velocity of the electromagnetic wave depends on the wavelength. In telecommunication the term of dispersion is used to describe the processes which cause that the signal carried by the electromagnetic wave and propagating in an optical fiber is degraded as a result of the dispersion phenomena. This degradation occurs because the different components of radiation having different frequencies propagate with different velocities. We distinguished various kinds of dispersion.

- As shown these phenomena are particularly important in optical telecommunication. In different periods of the historical development of the optical telecommunication the different kinds of dispersion played a different role. In the first period, when the multimode fibers were used and the light was transmitted only on small distances at low transmission speed, the chromatic dispersion played a negligible role in contrast to the mode dispersion. Then shown that the employing the single-mode optical fibers eliminated entirely the phenomenon of the mode dispersion and allowed to propagate the signal over large distances. However, with the higher transmission speeds gigabites per second the chromatic dispersion became more and more essential on large distances. When we apply a single-mode laser DFB (distributed feedback Bragg) of spectral width of 0.1 nm as a light source to propagate in a single-mode fiber characterized by the dispersion coefficient of 17 ps/nm/km, typical for most glassy fibers.

- Compensation necessary. However, dispersion compensation introduces additional loss and complexity to systems. This leads to various system configurations in terms of transmission fibers deployed. Each configuration has its own advantages and drawbacks. The system configuration is not the only key factor that improves system performance.
5.2 Recommendation

- The major parameters for the engineering of optical fiber communication networks using single mode fiber should be as follows:
  
  i. The mode structure.
  
i. Field distribution.
  
iii. Propagation constant.
  
iv. The management of losses (dB / km) and dispersion (ps / nm.km).

- In fiber optical high bit rate (such as 10G or 40G bit/s) long-haul transmission systems, dispersion compensation is one of the most important items to be considered for design.

- Management or optimization of residual dispersion are required for photonic networks, i.e., for fibers, repeaters and optical interfaces.

- PMD compensation is also required especially for 40Gbit/s or higher bit rate long-haul systems (PMD leads to broadening of optical pulses because of random variations in the birefringence of an optical fiber along its length. This broadening is in addition to GVD-induced pulse broadening. The use of dispersion management can eliminate GVD-induced broadening but does not affect the PMD induced broadening. For this reason, PMD has become a major source of concern for modern dispersion-managed systems.

- Some researchers (Mollenauer, 1991) are working on a novel method of propagating pulses in a fiber without dispersion. In the optical communications world there is immense interest in solitons [1]. This is due to the fact that solitons travel without dispersion of any kind on standard (highly dispersive) fiber. There is no need for dispersion compensation no matter how great the distance travelled. One of the consequences of this is that it becomes possible to have single-channel data streams of between 100 and 200 Gbps!\(^1\)

---

\(^1\)The word “soliton” is a contraction of the phrase “solitary solution” because the phenomenon represents a single solution to the propagation equation. Also, in physics, particles tend to be named with the suffix -on (electron, proton, neutron, photon...) and the word soliton therefore suggests particle-like behavior; chapter 9; Fiber-Optic Communications Systems, Third Edition. Govind P. Agrawal
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